

## Gas Absorption Studies in Microporous Hollow Fiber Membrane Modules

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A comprehensive experimental investigation of gas-liquid absorption in a shell-and-tube type microporous hydrophobic hollow fiber device in a parallel flow configuration was carried out. Two modes of countercurrent gas-liquid contacting were studied, the wetted mode (absorbent liquid filled pores) and the nonwetted mode (gas-filled pores). The absorbent flowed through the fiber bore in most of the experiments. The systems studied include pure CO<sub>2</sub>, pure SO<sub>2</sub>, CO<sub>2</sub>-N<sub>2</sub> mixtures and SO<sub>2</sub>-air mixtures. The absorbent was pure water. The absorption process was simulated for each case with a numerical model for species transport with and without chemical reaction. Laminar parabolic velocity profile was used for the tube-side flow, and Happel's free surface model was used to characterize the shell-side flow. The model simulations agreed well with the experimental observations in most cases. SO<sub>2</sub> removals as high as 99% were obtained in small compact contactors. High  $K_{La}$  and low height of transfer unit (HTU) values were obtained with hollow fiber contactors when compared to those of conventional contactors.

### Introduction

In conventional large-scale gas absorption processes, the gas mixture is contacted with a scrubbing liquor in devices like packed towers, spray towers, venturi scrubbers, and bubble columns. These dispersion-based contacting methods have many drawbacks. An approach having much promise and without these shortcomings is membrane-based contacting using microporous hydrophobic hollow fibers. The microporous membrane-based device acts as a gas absorber, with the gas flowing on one side and the absorbent liquid flowing on the other side of the membrane without either phase being dispersed in the other. The gas-liquid interface is at the pore mouth. Unlike that in a true liquid membrane process, the absorbent liquid has to be regenerated in a separate stripper which can be also a membrane-based contactor.

A membrane-based contacting device can have a number of configurations. The focus here will be primarily on parallel flow modules of microporous hollow fibers. The membranes can be hydrophobic or hydrophilic, and the pores can be either gas filled or absorbent filled. The absorbent may flow on the tube side or shell side. These membrane-based gas separation devices offer several advantages over conventional contacting devices: much higher surface area per unit contactor volume; independent control of gas and liquid flow rates without any flooding, loading, weeping, or foaming; known gas-liquid interfacial area; modularity, easy to scale up or down; availability of all of the membrane surface area for contacting regardless of how low the individual phase flow rates are.

These advantages have led a number of investigators to explore different applications. The first known application of a microporous membrane as a gas-liquid contacting device using hydrophobic flat Gore-Tex membranes of polytetrafluoroethylene was for oxygenation of blood (Esato and Eiseman, 1975). Tsuji et al. (1981) used hydrophobic microporous hollow fibers of polypropylene

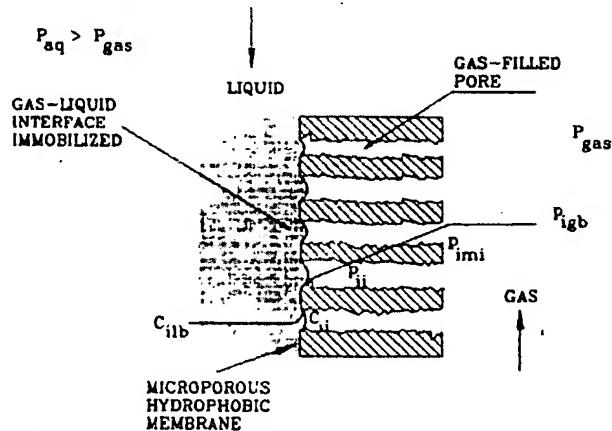


Figure 1. Nonwetted mode of hydrophobic membrane-based gas-liquid contacting.

for blood oxygenation. The applications of direct interest here are those for acid gas cleanup. Qi and Cussler (1985a,b) studied the absorption of a number of gases using various solvents and aqueous solutions in a hydrophobic microporous hollow fiber device where the membrane pores were gas filled. Yang and Cussler (1986) have noted a 10-fold increase in the gas separation efficiency when using a hollow fiber device over a conventional packed column. Ogundiran et al. (1988, 1989) in their experiments with absorption of SO<sub>2</sub> using microporous hydrophobic hollow fibers found this technique to be an extremely promising alternative to conventional scrubbers used for flue gas desulfurization. Sirkar (1992) has recently reviewed microporous membrane-based gas absorption processes.

The above-mentioned contacting studies were carried out using a microporous hydrophobic membrane. Figure 1 shows an absorbing liquid (that does not wet the fibers) flowing on one side and gas flowing on the other side. The gas pressure ( $P_{\text{gas}}$ ) has to be lower than that of the liquid ( $P_{\text{aq}}$ ) to prevent dispersion of gas as bubbles in the liquid. Liquid does not enter the pores unless a certain critical pressure is exceeded (Kim and Harriott, 1987). This mode of operation is identified as the nonwetted mode since the liquid does not wet the membrane and the membrane pores are gas filled. The gas/liquid interface shown in Figure

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## Experimental Section

Microporous hydrophilic Celgard X-20 and X-10 polypropylene hollow fibers made by Hoechst Celanese Corporation (SPD, Charlotte, NC) were used. Able to withstand air temperatures up to 75–80 °C, these fibers have an internal burst strength greater than 220 psi. Both fibers have a bubble point pressure of 150 psi and a mean pore diameter of 0.03 μm. The characteristics of hollow fiber modules 1–4 fabricated in the laboratory and used here are provided in Table I. Module 2 has the highest area/volume.

The gas composition was measured by a Hewlett-Packard 5890A gas chromatograph using a thermal conductivity detector (TCD). The detector response was recorded by a digital reporting integrator 3393A (H.P., Paramus, NJ). A 10-port gas sampling valve was utilized (Sengupta et al., 1990). A two-column system based on a 6-ft × 1/8-in. Chromosorb 108 column (Chrompack Inc., Bridgewater, NJ) and a 10-ft × 1/8-in. molecular sieve (Alltech Associates, Deerfield, IL) column was used. The column sequence was reversed after 1.9 min to separate O<sub>2</sub> and N<sub>2</sub> by the molecular sieve column. By this reversal CO<sub>2</sub>, SO<sub>2</sub>, and moisture did not enter the molecular sieve column thereby ensuring it a longer lifetime. Sampling was done continuously.

The liquid-phase concentrations were analyzed and used to compare the gas analysis via mass balances. In the case of pure CO<sub>2</sub> if the change in gas flow rate was limited, concentrations based on the liquid phase were considered more reliable. For CO<sub>2</sub>, liquid-phase-concentration measurements were made using a CO<sub>2</sub> electrode (Orion Model 94-02, Orion Research, Cambridge, MA). The measurement range of this electrode was 4.4–440 ppm CO<sub>2</sub> at a pH range of 4.8–5.2. Reproducibility was ±2%. Titration procedures were also used to cross-check the data (Betz, 1962). The concentrations measured by the electrode and the titration method were in agreement in the low range 0–100 ppm. In higher concentration ranges, a different titration method was used (Snell and Hitton, 1966). In the case of SO<sub>2</sub>, the method of analysis was based on the procedure described in *Standard Methods for the Examination of Water and Waste Water* (AWWA, 1965).

Figure 3 shows the experimental setup used for gas absorption studies. The gas mixture to be separated or the pure gas to be absorbed was passed through the module either on the shell side or on the tube side, and the absorbing liquid flowed countercurrently through the absorber on the other side, separated from the gas by the membrane walls. The gas flow rate was controlled by an electronic mass flow transducer and a flow controller (Models 8100 and 8200; Matheson, East Rutherford, NJ). Pressure gauges at the module inlet and outlet indicated the gas pressures which varied between 0 and 14 psig depending on the mode of operation. The gas stream outlet concentration was measured by a gas chromatograph (GC), after which the gas was vented into the laboratory exhaust hood. A needle valve was used to control the liquid flow rate. The liquid flow rate was measured at the outlet of the membrane contactor at regular intervals. The back-pressure regulator (BPR) was not needed when the membrane pores were gas filled; i.e., the system was operated in the unwetted mode. Experimentally observed concentration change from the feed to the outlet of the module in the gas phase was used to determine the overall logarithmic mean based overall mass-transfer coefficient  $K_{Olm}$ , using the following equation:

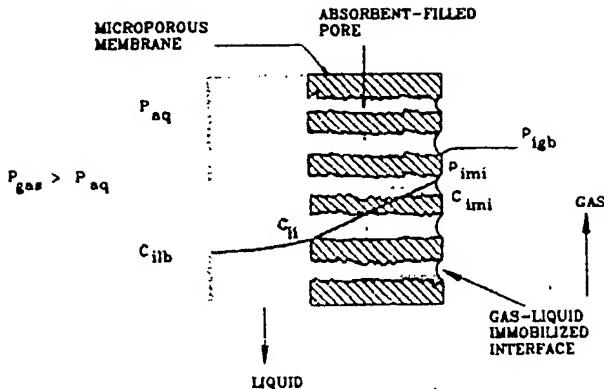


Figure 2. Wetted mode of membrane-based gas–liquid contacting.

1 lies at the pore mouth of the hydrophobic fiber on the absorbing liquid side. In this mode the gas species being separated diffuses first through the gas in the membrane pores before dissolving in the absorbent. Because the pores in the fibers are gas filled, the gas-phase mass-transfer coefficients are much higher than the liquid-side mass-transfer coefficients if no fast reaction is taking place or when the gas solubility is very low. Thus gas absorption is likely to occur without practically any increase in mass-transfer resistance due to the presence of the membrane.

There is yet another mode of operation where the membrane pores are absorbent liquid filled; this is called the wetted mode and is shown in Figure 2. Here a hydrophilic membrane is spontaneously wetted by the absorbing liquid which flows on one side, while the gas flows on the other side. In the case shown in Figure 2, the gas pressure has to be higher than the liquid pressure to prevent the liquid from dispersing as drops in the gas. A hydrophobic membrane can also be used in this mode using a nonwetting liquid by employing an exchange process (Bhave and Sirkar, 1986, 1987) which allows one to incorporate a nonwetting absorbing solution in the hydrophobic membrane pores. This mode may be useful when there is a very fast or instantaneous liquid-phase reaction and the gas resistance would be controlling the rate of transfer. The membrane-phase resistance may be eliminated in such a mode of operation. There is another reason for investigating such a mode. If long-term operation of a hydrophobic membrane leads to wetting of the pores, such a mode of operation may be adopted.

In this paper, the absorption of acid gases like SO<sub>2</sub> and CO<sub>2</sub> have been investigated using modules of microporous hydrophobic hollow fine fibers. Experiments have been carried out with both gas-filled membrane pores (non-wetted mode) and liquid-filled membrane pores (wetted mode) using pure gases and their mixtures, with different absorbing liquids flowing either on the tube side or on the shell side. This paper focuses on pure water as absorbent. A second paper will cover diethanolamine in water. Although the absorption of CO<sub>2</sub> in water is essentially physical, absorption of SO<sub>2</sub> in water is very strongly influenced by ionization so that there is reactive absorption. Gas absorption models have been developed on the basis of first principles to explain the observed separation and evaluate some commonly used correlations. The present analysis attempts to provide a fundamental basis for understanding aspects of this promising contacting technique. The experimental values of  $K_{1,0}$  and HTU obtained here have been compared with those available in the literature. Gaseous systems investigated here include pure CO<sub>2</sub>, pure SO<sub>2</sub>, CO<sub>2</sub>-N<sub>2</sub> mixtures, and SO<sub>2</sub>-air mixtures.

Table I. Characteristics of Hollow Fiber Modules Used in This Study

| module no. | fiber o.d., $\mu\text{m}$ | fiber i.d., $\mu\text{m}$ | shell diam., cm | length of module, cm | no. of fibers | porosity | packing fraction | area of contact i.d., $\text{cm}^2$ | $a, \text{cm}^2/\text{cm}^3$ |
|------------|---------------------------|---------------------------|-----------------|----------------------|---------------|----------|------------------|-------------------------------------|------------------------------|
| 1          | 150                       | 100                       | 0.4572          | 30.0                 | 102           | 0.2      | 0.1098           | 96.13                               | 19.52                        |
| 2          | 300                       | 240                       | 0.4927          | 30.0                 | 132           | 0.4      | 0.4887           | 298.57                              | 52.2                         |
| 3          | 300                       | 240                       | 0.4572          | 18.0                 | 75            | 0.4      | 0.3229           | 101.78                              | 34.44                        |
| 4          | 300                       | 240                       | 0.4572          | 17.5                 | 75            | 0.4      | 0.1098           | 98.96                               | 34.44                        |

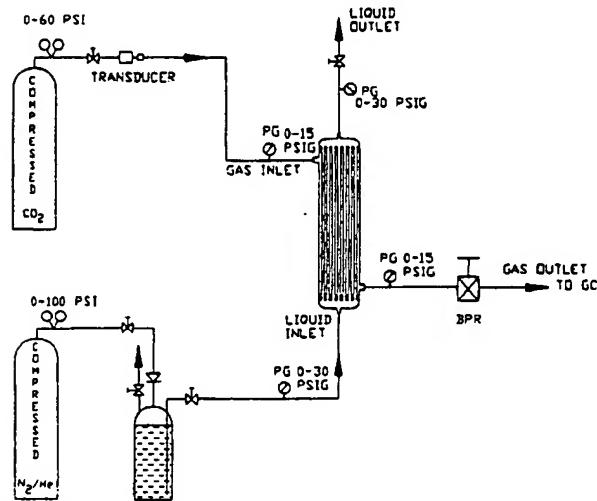


Figure 3. Gas absorption using hollow fibers: experimental setup.

$$K_{0lm} = \frac{V_l C_{l,out}}{60 A_T \Delta C_{lm}} \quad (1)$$

$K_{0lm}$  is also denoted as  $K_L$  in this paper.  $A_T$  is based on the surface area of gas-liquid contact. The quantity  $\Delta C_{lm}$ , the logarithmic mean concentration difference of gas species, is given by

$$\Delta C_{lm} = \frac{(HC_{g,in} - C_{l,out}) - (HC_{g,out})}{\ln\left(\frac{HC_{g,in} - C_{l,out}}{HC_{g,out}}\right)} \quad (2)$$

Here  $H$  is Henry's law constant which is the ratio of the interfacial concentration of absorbed species  $i$  in liquid divided by that in the gas:

$$C_{il} = HC_{ig} = C_{Aw} \quad (3)$$

if species  $i$  is represented by  $A$  and the liquid phase by  $w$  for water. In the case of  $\text{CO}_2\text{-H}_2\text{O}$  system, the value of Henry's law constant is 0.8314 at 25 °C (Geankoplis, 1972) and that for  $\text{SO}_2\text{-H}_2\text{O}$  is 25.86 at 25 °C (Hikita et al., 1978).

### Model Development

A theoretical model is developed here to describe the absorption of pure  $\text{CO}_2$ , pure  $\text{SO}_2$ , and their mixtures in nitrogen or air by water. Both wetted and nonwetted modes of operation are considered here, and in both modes, the liquid flow was always countercurrent to the gas flow and on the tube side. The following assumptions were utilized: (1) steady state and isothermal condition; (2) no axial diffusion; (3) fully developed parabolic liquid velocity profile in the tube side; (4) use of Happel's free surface model (1959) to characterize the shell-side velocity profile; (5) ideal gas behavior; (6) applicability of Henry's law.

Using these assumptions, the general form of the conservation equation for the absorption processes is given as

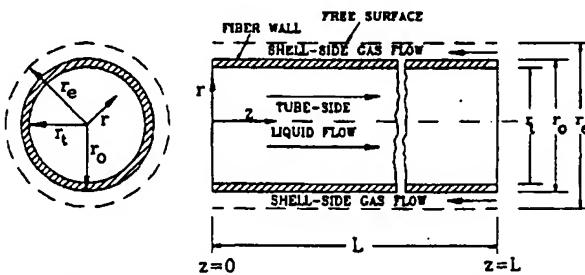


Figure 4. Free surface model (adapted from Gill and Bansal (1973)).

$$D_{ik} \left[ \frac{\partial^2 C_{ijk}}{\partial r^2} + \frac{1}{r} \frac{\partial C_{ijk}}{\partial r} \right] = v_{zh}(r) \frac{\partial C_{ijk}}{\partial z} + (-r_i) \quad (4)$$

In the above equation,  $i$  denotes the components involved ( $i = A, B, \dots$ ),  $j$  denotes the phase (gas (g), liquid (l)), and  $k$  denotes the location (shell (s), tube (t), or membrane (m)). The quantity  $v_{zh}$  is the prescribed axial velocity profile for the tube or the shell side. For the tube side a laminar parabolic velocity profile is assumed:

$$v_{zt} = 2(v_{zz})[1 - (r/r_t)^2] \quad (5)$$

The flow in the shell side is usually very complex. Happel's (1959) free surface model as represented in Figure 4 is used to characterize the shell side velocity profile; it is given by the following equation:

$$v_{zs}(r) = 2(v_{zz}) \left[ 1.0 - \left( \frac{r_0}{r_s} \right)^2 \right] \times \left[ \frac{(r/r_s)^2 - (r_0/r_s)^2 + 2 \ln(r_0/r)}{3 + (r_0/r_s)^4 - 4(r_0/r_s)^2 + 4 \ln(r_0/r_s)} \right] \quad (6)$$

where  $r_s$  is the radius of the free surface and is given as

$$r_s = \left[ \frac{1}{1-\epsilon} \right]^{1/2} r_0$$

**Pure Gas Absorption: Nonwetted mode.** In this mode of absorption of pure  $\text{CO}_2$ , the pores of the membrane are filled with pure gas flowing in the shell side. There is no gas-phase-transport resistance. Due to the low value of the equilibrium constant ( $K_E = 4.4 \times 10^{-7} \text{ mol/L}$ ), bicarbonate formation is neglected. Equations 4 and 5 are solved for the liquid phase for  $i = A = \text{CO}_2$ ,  $j = l$ ,  $k = t$ , and  $r_A = 0$ , along with the following boundary conditions. The concentration profile is symmetrical about the center of the tube and therefore satisfies

$$\frac{\partial C_{ilt}}{\partial r}(0,z) = 0 \quad (7)$$

Further, since pure  $\text{CO}_2$  gas of known pressure is present in the shell side, the species concentration at the inside wall ( $r = r_t$ ) is known. Hence the boundary condition at the wall is

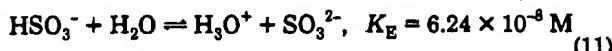
$$C_{ilt}(r_t, z) = C_{iw} \quad (8)$$

Here  $C_{iw}$  is the product of Henry's law constant times gas-phase concentration and is given by eq 3. In most of the experiments conducted, the pure gas pressure remained

essentially constant along the module length. Therefore  $C_{iw}$  remained constant with  $z$  (which is essentially the Graetz solution (Skelland, 1974)). There were a few experimental situations where the pure gas pressure changed somewhat along the module. In such a case,  $C_{iw}$  varied with axial location  $z$  (linear variation was assumed for numerical calculations). The inlet condition is also known

$$C_{it}(r_0, 0) = C_{inlet} = 0 \quad (9)$$

So far in the species conservation equation, there was no reaction. The situation changes for absorption of pure  $\text{SO}_2$  in water. The following hydrolysis reactions take place when  $\text{SO}_2$  is absorbed in pure water:



If the pH is less than 4–5, formation of  $\text{SO}_3^{2-}$  is negligible. By choosing operation conditions that exclude the pH range mentioned above,  $\text{SO}_3^{2-}$  formation is avoided (Seinfeld, 1980; Roberts, 1979). In the present study the operational conditions were such as to allow the assumption that only  $\text{HSO}_3^-$  formation is effective.

The diffusion coefficients of  $\text{HSO}_3^-$  (species B) and  $\text{H}_3\text{O}^+$  at 25 °C are  $1.47 \times 10^{-5}$  and  $9.3 \times 10^{-5} \text{ cm}^2/\text{s}$ , respectively. An electrical potential develops which retards the motion of  $\text{H}_3\text{O}^+$  and speeds up  $\text{HSO}_3^-$ , so as to maintain electroneutrality. Thus an effective diffusion coefficient in the liquid is defined, which in the binary case is given as (Vinograd and McBain, 1941)

$$D_{Bi} = \frac{2D_{\text{H}_3\text{O}^+}D_{\text{HSO}_3^-}}{D_{\text{H}_3\text{O}^+} + D_{\text{HSO}_3^-}} \quad (12)$$

The rate expressions for the reaction (eq 10), after using the electroneutrality condition,  $C_{\text{HSO}_3^-} = C_{\text{H}_3\text{O}^+}$ , are expressed by the following equations:

$$-r_{\text{SO}_2} = k_1 C_{\text{SO}_2} - \frac{k_1}{K_E} (C_{\text{HSO}_3^-})^2 \quad (13)$$

$$-r_{\text{HSO}_3^-} = -k_1 C_{\text{SO}_2} + \frac{k_1}{K_E} (C_{\text{HSO}_3^-})^2 \quad (14)$$

where the forward rate constant  $k_1$  is given as  $3.17 \times 10^{-2} \text{ s}^{-1}$  (Wang and Himmelblau, 1964). Since eq 4 is valid for both components  $\text{SO}_2$  (A), and  $\text{HSO}_3^-$  (B), there are now two coupled partial differential equations to be solved along with the boundary conditions. With eqs 13 and 14 representing the reaction terms  $-r_A$  and  $-r_B$ , respectively, eq 4 for  $\text{SO}_2$  and eq 4 for  $\text{HSO}_3^-$  are solved simultaneously along with the above boundary conditions (eqs 7 and 8). In the case of  $\text{HSO}_3^-$  at the wall, equilibrium is assumed and the following equation is used to represent the wall condition:

$$C_{Bl}(r_0, z) = (C_{At} K_E)^{1/2} \quad (15)$$

For details, see the thesis by Karoor (1992).

**Wetted Mode.** In this mode of operation the liquid is flowing in the tube and is also incorporated in the membrane pores. An additional equation is required to describe the diffusion through the membrane pores as the liquid-filled pores offer additional resistance. At high porosity (greater than 10%), Keller and Stein (1967) have shown that the diffusional process is essentially one

dimensional. The following assumptions were made in deriving the membrane-phase equation:

1. The effective membrane-phase diffusion coefficient  $D_{im}$  is defined by  $D_{it}\epsilon/\tau$ .
2. The diameter of the pore is uniform.
3. There are no gradients in the direction normal to the pore axis.
4. There is no convective transport.

The species balance eq 4 is still valid for the liquid in the membrane phase ( $j = l, k = m$ ) for each species without the convection term. The following boundary conditions are used at the membrane-tube interface:

$$r = r_t, \quad C_{ilm}|_{r_t} = C_{ilt}|_{r_t} \quad (16)$$

A flux continuity at the same boundary is assumed, giving the equation

$$r = r_t, \quad D_{im} \frac{\partial C_{ilm}}{\partial r}|_{r_t} = D_{lt} \frac{\partial C_{ilt}}{\partial r}|_{r_t} \quad (17)$$

The equation for the membrane phase is solved using the above boundary conditions, namely eqs 16 and 17, to yield the following equation at the inner wall of the membrane  $r = r_t$  (Karoor, 1992).

$$\frac{\partial C_{At}}{\partial r}|_{r_t} = [C_{Alm}|_{r_0} - C_{Alt}|_{r_t}] \frac{D_{Am}}{D_{At}} \frac{1}{(\ln r_0 - \ln r_t)r_t} \quad (18)$$

Now eq 4 for  $i = A, j = l$ , and  $k = t$  along with eq 18, the boundary condition eq 7, and the wall condition at  $r = r_0$ ,  $C_{ilm}(r_0, z) = C_{iw}$  are solved.

The above system of equations is for  $\text{CO}_2$ , since there is no reaction in the liquid phase. For  $\text{SO}_2$ , two species  $\text{SO}_2$  (A) and  $\text{HSO}_3^-$  (B) are present and eq 4 is valid for each species for the tube side. The equations for the membrane phase are given as

$$D_{Am} \left[ \frac{\partial^2 C_{Alm}}{\partial r^2} + \frac{1}{r} \frac{\partial C_{Alm}}{\partial r} \right] = k_1 C_{Alm} - \frac{k_1}{K_E} C_{Blm}^2 \quad (19)$$

$$D_{Bm} \left[ \frac{\partial^2 C_{Blm}}{\partial r^2} + \frac{1}{r} \frac{\partial C_{Blm}}{\partial r} \right] = -k_1 C_{Alm} + \frac{k_1}{K_E} C_{Blm}^2 \quad (20)$$

A flux continuity equation at the boundary is assumed, giving the equations

$$r = r_t, \quad D_{Am} \frac{\partial C_{Alm}}{\partial r}|_{r_t} = D_{At} \frac{\partial C_{Alt}}{\partial r}|_{r_t} \quad (21)$$

$$r = r_t, \quad D_{Bm} \frac{\partial C_{Blm}}{\partial r}|_{r_t} = D_{Bt} \frac{\partial C_{Blt}}{\partial r}|_{r_t} \quad (22)$$

Equations 19 and 20 are solved along with the boundary conditions eqs 21 and 22 and then combined with the equations for the interior and the boundary conditions.

**Absorption from a Gas Mixture.** In gas absorption from a mixture, the gas-phase resistance has to be considered; an additional equation for the gas phase has to be solved. For  $\text{CO}_2$  absorption in water, eq 4 is also valid for the gas flowing in the shell side. The shell-side gas velocity profile is given by eq 6. At the free surface  $r = r_0$ , the following boundary condition is valid:

$$\frac{\partial C_{ig}}{\partial r}(r_0, z) = 0 \quad (23)$$

At two membrane interfaces the following conditions are valid:

$$r = r_0, \quad C_{igm}|_{r_0} = \frac{C_{il}|_{r_0}}{H} \quad (24)$$

Table II. Parameters Used in the Absorption of SO<sub>2</sub> and CO<sub>2</sub> in Water

| temp<br>°C | solubility <sup>a</sup><br>of<br>SO <sub>2</sub> ( $H_{iw}$ ),<br>mol/liter | diffusivity <sup>a</sup><br>of SO <sub>2</sub> ( $10^6 D$ ),<br>cm <sup>2</sup> /s | solubility <sup>b</sup><br>of CO <sub>2</sub> ( $H_{iw}$ ),<br>mol/L | diffusivity <sup>c</sup><br>of CO <sub>2</sub> ( $10^6 D$ ),<br>cm <sup>2</sup> /s | equilib <sup>a</sup> constant<br>of SO <sub>2</sub> ( $K$ ),<br>mol/L |
|------------|---|--|--|--|---|
| 15         | 1.5420  | 1.34   | 0.045448   | 1.40   | 2.19  |
| 25         | 1.0559  | 1.76   | 0.3388   | 1.92   | 1.70  |

<sup>a</sup> Hikita et al. (1978). <sup>b</sup> Geankoplis (1972). <sup>c</sup> Majumdar et al. (1988).

$$r = r_0, \quad C_{igm}|_{r_0} = C_{igl}|_{r_0} \quad (25)$$

Flux continuity equations at the same boundaries are also valid, giving the equations

$$r = r_t, \quad D_{im} \frac{\partial C_{igm}}{\partial r}|_{r_t} = D_{it} \frac{\partial C_{ilt}}{\partial r}|_{r_t} \quad (26)$$

$$r = r_0, \quad D_{im} \frac{\partial C_{igm}}{\partial r}|_{r_0} = D_{is} \frac{\partial C_{igs}}{\partial r}|_{r_0} \quad (27)$$

Then, eq 4 is solved for the membrane phase with the above boundary conditions to give the gradients at  $r = r_t$  and  $r = r_0$ , as

$$\frac{\partial C_{Alt}}{\partial r}|_{r_t} = \left[ C_{Ags}|_{r_0} - \frac{C_{Alt}|_{r_t}}{H} \right] \frac{D_{Am}}{D_{At}} \frac{1}{[\ln r_0 - \ln r_t] r_t} \quad (28)$$

$$\frac{\partial C_{Ags}}{\partial r}|_{r_0} = \left[ C_{Ags}|_{r_0} - \frac{C_{Alt}|_{r_t}}{H} \right] \frac{D_{Am}}{D_{As}} \frac{1}{[\ln r_0 - \ln r_t] r_0} \quad (29)$$

The equations for the shell and tube are combined with the above equations for the membrane inner and outer boundaries and solved simultaneously with the boundary conditions eqs 7 and 23 and the inlet conditions for CO<sub>2</sub> absorption. In the case of SO<sub>2</sub> absorption all the above equations and boundary conditions are valid except that the reaction term has to be taken into account; in addition similar relations are written for HSO<sub>3</sub><sup>-</sup>. At the gas-liquid interface the formation of HSO<sub>3</sub><sup>-</sup> is assumed to be zero (for details, see Karoor (1992)).

The above equations are solved using a finite difference scheme (Davis, 1984). The diffusion terms are discretized using the centered in space scheme, whereas the convection term is discretized using the upwind difference scheme. A downward difference scheme is used when the velocity in the shell side is in the opposite direction to that of the tube side. Using the method of lines, the above partial differential equations were reduced to a set of algebraic equations. These equations are tridiagonal in form as a result of the discretization method adopted. This set can now be easily solved using the Thomas algorithm (Karoor, 1992). In the case of SO<sub>2</sub> absorption the equation is solved for  $C_{Alt}$  (based on  $C_{Blt}$  values determined in the earlier iteration). The equations for  $C_{Blt}$  are solved using a similar approach; the nonlinear term in eqs 19 and 20 are pseudolinearized by representing  $C_{Blm}^2$  and  $C_{Blm} C_{Blm}^*$ , where  $C_{Blm}^*$  is the value of  $C_{Blm}$  from the previous iteration. The solution procedure is continued till convergence is attained.

## Results and Discussion

The results of countercurrent absorption of pure CO<sub>2</sub> and pure SO<sub>2</sub> in water flowing in the tube side of the hollow fiber modules are presented and discussed first. Performance of gas mixtures is provided later. The values of the parameters used are given in Table II. The experimental results have been compared with the results obtained from the numerical solution of the governing partial differential equations. Limited performance data have been provided also for shell-side water flow. Both experimental and

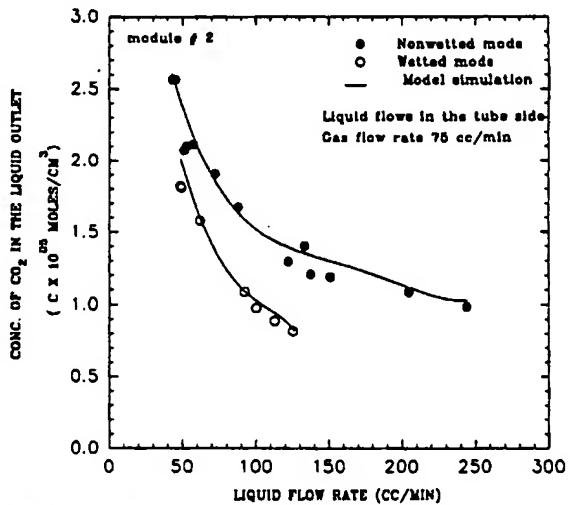


Figure 5. Comparison between wetted and nonwetted modes of operation for absorption of pure CO<sub>2</sub> in water.

model results have also been recast in the fashion of a plot of Sherwood number vs Graetz number. Besides the pure gas absorption experiments, the experimental results and the model simulations are discussed for cases where the feed gas is either a mixture of CO<sub>2</sub> (40%, 10%) in N<sub>2</sub> or a mixture of SO<sub>2</sub> (1%, 0.95%, 0.4%) in air. One or two or all three of the hollow fiber modules 2, 3, and 4 were used in these studies.

Figure 5 illustrates the variation of the liquid outlet concentration of CO<sub>2</sub> as a function of water flow rate in module 2. The gas flow rate was held constant at 75 cm<sup>3</sup>/min for all the runs shown in the figure as the liquid flow rate was varied. Data and numerical simulation results are presented for both nonwetted mode and wetted mode. It is clear that the liquid outlet CO<sub>2</sub> concentration in the wetted mode is lower than that in the nonwetted mode. This effect is more apparent as the liquid velocity is increased. Since the liquid boundary layer resistance decreases with increasing liquid velocity, the membrane resistance starts to control. By filling the pores with the absorbent liquid, the membrane-phase resistance is increased as the liquid offers additional resistance to the transport; this behavior highlights the fact that this is a liquid-phase-controlled process. Hence it would not be advantageous to operate under wetted-mode conditions for this particular system. Exactly the same trends were observed in the case of module 3. The theoretical curves shown in Figure 5 were obtained by solving eq 4 along with the boundary conditions. The theoretical predictions for these situations have a maximum deviation of 7% from the experimental results.

The effect of using a module with a higher contact area can be observed by using a greater number of fibers or increasing the length of the module. Figure 6 shows a plot of the concentration of CO<sub>2</sub> in the liquid outlet for three different modules of varying surface area as a function of the liquid flow rate in the nonwetted mode. From Table I it is seen that module 2 has the maximum contact area followed by module 3 and then module 4. In the case of

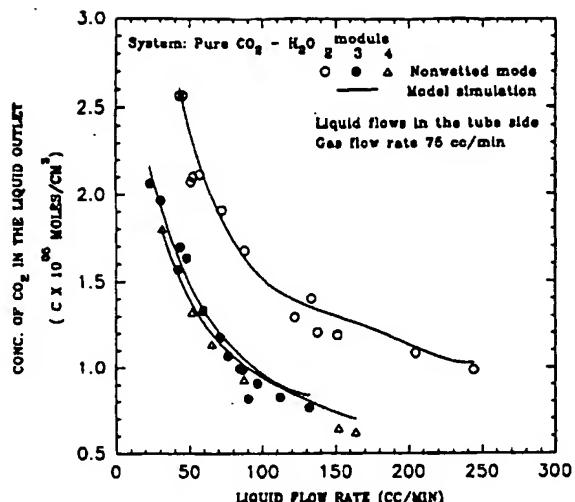


Figure 6. Comparison of concentration of  $\text{CO}_2$  in the liquid outlet for modules of different surface area.

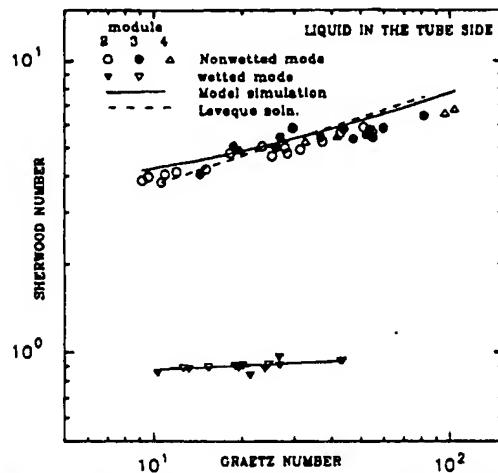


Figure 7. Sherwood number versus Graetz number variations for pure  $\text{CO}_2$  absorption in water with different modules.

module 2, a higher amount of  $\text{CO}_2$  in the liquid outlet is observed for the same liquid flow rate when compared with modules 3 and 4. Obviously using a longer module would require a higher pressure drop to achieve a given flow velocity. A maximum  $\text{CO}_2$  removal of 76% was achieved in module 2 with a liquid flow rate of 244  $\text{cm}^3/\text{min}$ . The numerical model simulations describe these behavior quite well.

Figure 7 illustrates the same experimental results for modules 2, 3, and 4, but now as a plot of Sherwood number versus Graetz number, where Sherwood number is defined as  $K_{\text{Obm}}(2r_t)/D_{\text{At}}$  for the nonwetted mode (for the wetted mode  $r_t$  is replaced by  $r_0$ ) and Graetz number for the tube side is defined as  $N_{\text{Re}}N_{\text{Sc}}\pi r_t/2L$ .  $K_{\text{Obm}}$  is the logarithmic mean overall mass-transfer coefficient calculated by eq 1. A logarithmic mean driving force is used in the calculations as given by eq 2. An increase in Graetz number due to an increase in the liquid flow rate increases the mass-transfer coefficient as expected. Another feature that can be observed is the effect of operating the module in the wetted mode. In the wetted mode the pores of the fiber are filled with the absorbent liquid. This increases the resistance to mass transfer across the membrane pressure, as a result of decreased gas diffusion coefficient in the liquid phase. This effect is clearly shown in Figure 7 via a decrease in the Sherwood number by almost an order of

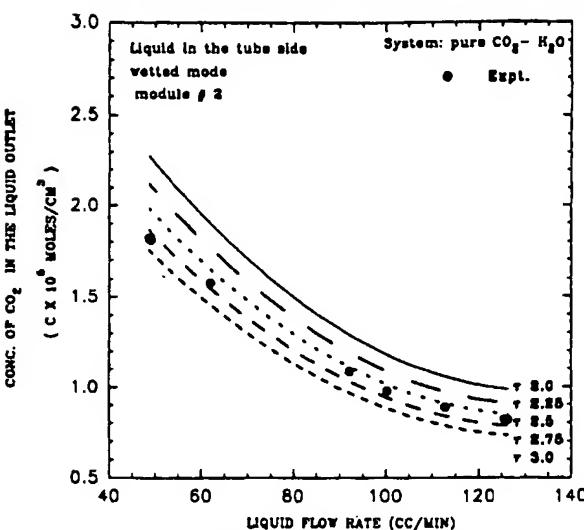


Figure 8. Effect of tortuosity on the  $\text{CO}_2$  outlet concentration for module 2.

magnitude. As is evident from Figure 7, the effect of the wetted mode in module 3 was also observed to be similar. The extent of decrease in mass transfer in the wetted mode in module 3 is also comparable to that of module 2.

Figure 7 also illustrates the results obtained from the theoretical model for the nonwetted mode. The model results and the experimental observations show good agreement with a maximum deviation of about 11.5%. The Leveque solution (Leveque, 1928), also plotted in this figure, shows a reasonable agreement with the data. The Leveque solution results are somewhat higher at higher Graetz numbers.

In the wetted mode, an effective membrane diffusion coefficient was defined as  $D_{\text{Am}} = D_{\text{At}}/\tau$ , where  $\epsilon/\tau$  was taken to be 0.4/2.5, from the experimentally measured values by Prasad and Sirkar (1988) for X-20 fibers with 240- $\mu\text{m}$  inner diameter and 300- $\mu\text{m}$  outer diameter. In the case of X-10 hollow fibers with 100- $\mu\text{m}$  inside diameter and 150- $\mu\text{m}$  outside diameter, the quantity  $\epsilon/\tau$  was taken to be 0.2/3.0.

Using this definition of effective diffusion coefficient and the values of  $\epsilon/\tau$ , the model discussed before for the wetted mode was solved numerically to obtain the results shown in Figure 7. The agreement between the model and experimental results is good with a maximum deviation of 7.5% for module 2. For module 3, Figure 7 shows that the agreement between the model results and experiments is again good with a maximum deviation of 12%. As discussed before, the length of module 2 is higher than that of module 3; also the number of fibers is greater in module 2 when compared to module 3. Thus, under similar operating flow rates, the Graetz number is lower for module 2 when compared to module 3, and it results in lower Sherwood numbers for module 2.

In Figure 8 the predicted  $\text{CO}_2$  outlet concentration for module 2 in the wetted mode is plotted by allowing a maximum variation of  $\pm 20\%$  in the tortuosity  $\tau$  for X-20 fibers (2.5) used in the present study. As the tortuosity increases, the transfer of  $\text{CO}_2$  into the liquid decreases since  $\text{CO}_2$  has to diffuse through longer lengths of liquid. If the measured tortuosity is  $\pm 20\%$  different from the value 2.5 used, then the maximum difference in the concentration is around 18% for the liquid flow rates and the system under study.

Figure 9 illustrates the results for module 3 when pure  $\text{SO}_2$  flows in the shell side and water in the tube side. The

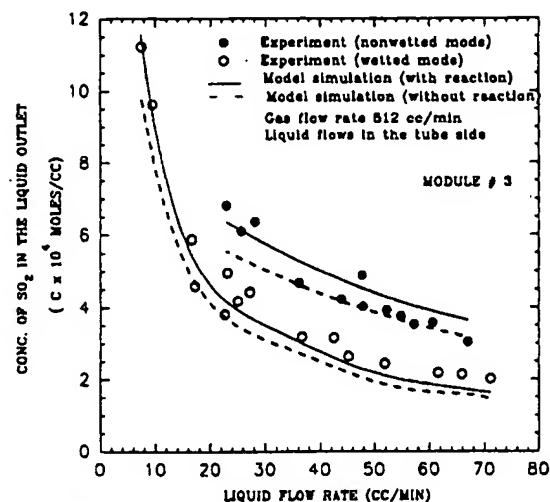


Figure 9. Absorption of pure  $\text{SO}_2$  in water: concentration variation as a function of liquid flow rate for module 3.

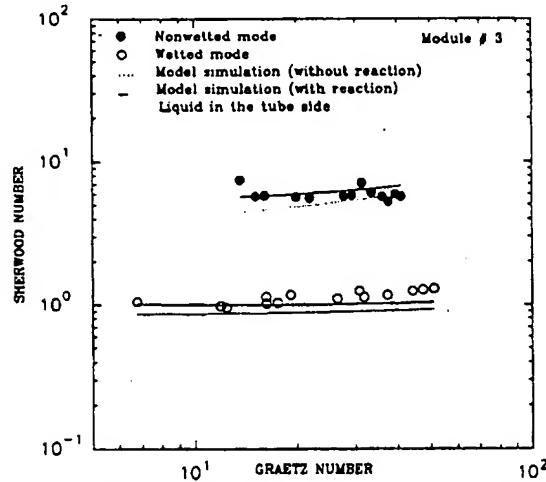


Figure 10. Sherwood number versus Graetz number for absorption of pure  $\text{SO}_2$  in water.

experimental results have been indicated for both wetted and nonwetted modes; the wetted mode again offers higher resistance to mass transfer when compared to the nonwetted mode. Even though the solubility of  $\text{SO}_2$  in water is high and there is a dissociation reaction, the reaction is not fast enough to make the liquid-filled-membrane resistance negligible. The theoretical predictions obtained by numerically solving the coupled differential equations explain the experimental results well with a maximum deviation of about 13%. The equations were also solved assuming that there is no reaction occurring between  $\text{SO}_2$  and  $\text{H}_2\text{O}$ . It was found that there was a 10%–15% increase in predicted  $\text{SO}_2$  concentration at the liquid outlet when the reaction term was included. Generally the latter approach describes the data much better especially in the wetted mode. It is worthwhile to note that a more rigorous analysis of  $\text{SO}_2$  dissociation in water could have been carried out in view of the reaction being an instantaneous reversible one. However, the scatter in the data (Figure 9, nonwetted case) is considerable, and the enhancement due to  $\text{SO}_2$  ionization in water is known to be around 10% (Danckwerts, 1970).

Figure 10 is a plot of Sherwood number versus Graetz number for the absorption of pure  $\text{SO}_2$  in water. The variation of Sherwood number with Graetz number is not very high in this case, due to the very high  $\text{SO}_2$ 's solubility

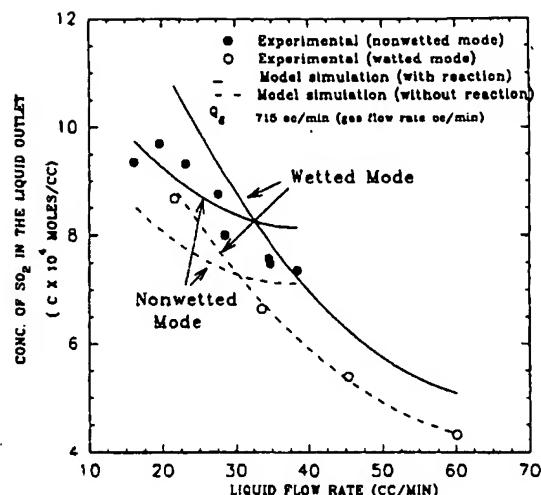


Figure 11. Absorption of pure  $\text{SO}_2$  in water: liquid flowing in the tube side of module 2.

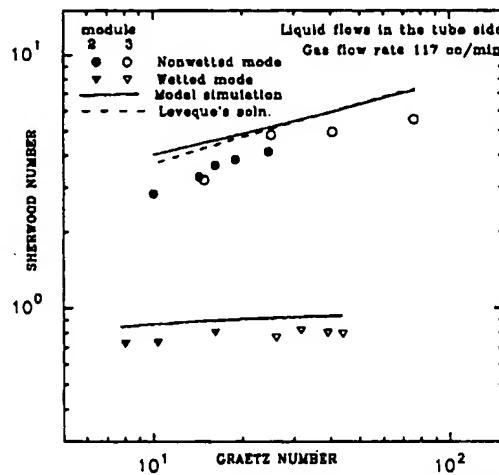


Figure 12. Sherwood number versus Graetz number for 10%  $\text{CO}_2$  feed gas absorption in water.

in water and the reaction that is occurring in the liquid phase. With a system having higher solubility or faster reaction rates, we may see a transition to gas-phase control (see Karoor (1992) for  $\text{SO}_2$ - $\text{Na}_2\text{SO}_3$  system).

Figure 11 shows the results for the same configuration with module 2 for pure  $\text{SO}_2$  and water. At low liquid flow rates, when wetted and nonwetted modes of operations are considered, the difference in the concentration of  $\text{SO}_2$  at the liquid outlet is small since the gas-liquid contact area is high and the liquid flow rate is low so that the liquid gets saturated before passing the whole length of the module.

So far, the cases studied utilized pure gas in the shell side; therefore the gas-phase resistance was negligible. Figure 12 shows the results for modules 2 and 3 when a feed gas containing 10%  $\text{CO}_2$  in  $\text{N}_2$  was flowing in the shell side and liquid was flowing countercurrently in the tube side. The results are presented as a plot of Sherwood number versus Graetz number. An increased resistance to mass transfer in the wetted mode is again apparent. The results of the model simulations show a maximum deviation of about 25% in the case of nonwetted mode. In the case of the wetted mode, the model seems to explain the experimental observations within a maximum deviation of 18%. These deviations may have resulted from all of the gas not coming into contact with the liquid in

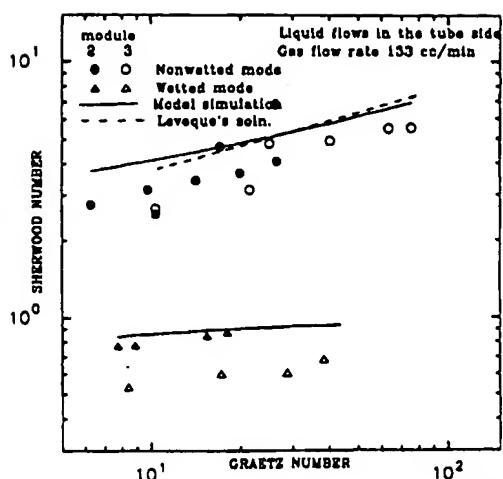


Figure 13. Sherwood number versus Graetz number for 40%  $\text{CO}_2$  feed gas absorption in water.

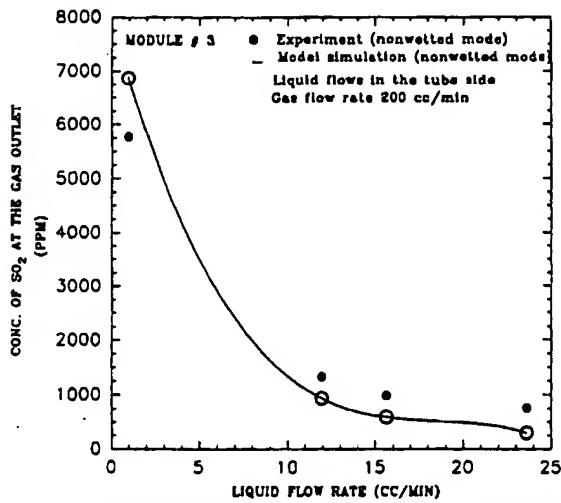


Figure 14.  $\text{SO}_2$  absorption in water from 1%  $\text{SO}_2$ -air mixture in module 3.

the modules which do not have high packing density. Similar trends were obtained in the case of the 40%  $\text{CO}_2$ -balance  $\text{N}_2$  gas mixture (Figure 13); the maximum deviation in the nonwetted mode was around 25% in the case of module 2 and around 25% in the case of module 3. In the wetted mode, the deviation was 10% for module 2 and 35% in the case of module 3. This may be attributed to bypassing in the shell-side flow and to the fact that the packing fraction is lower in the case of module 3 when compared to module 2.

In these calculations the diffusion coefficient of the gas in the pores in the nonwetted mode was calculated from the transition regime of gas diffusion in porous media. Since  $r_p/\lambda$  was around 0.56, the diffusion of  $\text{CO}_2-\text{N}_2$  in the pore was considered to be in the transition regime and not in the Knudsen regime. It was found that if the molecular diffusion coefficient was used there was no difference in the results as  $\text{CO}_2$  absorption in water is liquid phase controlled.

Figure 14 identifies the concentration of  $\text{SO}_2$  at the gas outlet for module 3 when 1%  $\text{SO}_2$  (balance air) mixture was scrubbed using water in module 3. It is seen that, as the liquid flow rate increases, the concentration of  $\text{SO}_2$  decreases, as expected. When the liquid flow rate was increased above 40  $\text{cm}^3/\text{min}$ , essentially no  $\text{SO}_2$  was observed in the gas outlet. For module 2 at the same gas

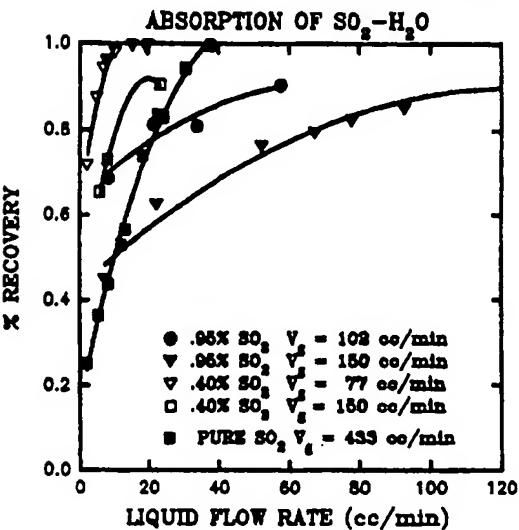


Figure 15. Absorption of  $\text{SO}_2$  from feed gases of pure  $\text{SO}_2$  and  $\text{SO}_2$ -containing gas mixtures in water flowing in the shell side.

flow rate ( $200 \text{ cm}^3/\text{min}$ ), no  $\text{SO}_2$  was observed experimentally at the gas outlet even at much lower liquid flow rates. This is because of the higher contacting area in the case of module 2 when compared with module 3. Thus it is observed that, with such small modules, it was possible to absorb all of the  $\text{SO}_2$  using simple water as the absorbent.

For the systems  $\text{CO}_2-\text{H}_2\text{O}$  and  $\text{SO}_2-\text{H}_2\text{O}$  studied, the rate of absorption of the gas species scrubbed using the same module is much higher for  $\text{SO}_2$  when compared to  $\text{CO}_2$ . There are several reasons for this: the solubility of  $\text{SO}_2$  in water is 2 orders of magnitude higher than that of  $\text{CO}_2$  and also the dissolution of  $\text{SO}_2$  in water is accompanied by a reaction. Even though  $\text{CO}_2$  reacts with water, the rate is orders of magnitude slower when compared to  $\text{SO}_2$ . Thus the absorption of  $\text{SO}_2$  by water is enhanced by both the increased solubility and the chemical reaction.

It is useful now to identify the level of gas purification achieved using small hollow fibers modules having high surface areas. Usually most of the  $\text{CO}_2$  and  $\text{SO}_2$  were easily scrubbed out. In the case of  $\text{SO}_2$ , removal as high as 98% was achieved with a liquid flow rate of  $38.46 \text{ cm}^3/\text{min}$  and gas flow rate of  $714.7 \text{ cm}^3/\text{min}$  in the 30-cm-long, 0.49-cm-diameter module 2 from a pure  $\text{SO}_2$  feed gas. When water is used as an absorbent, it is better to operate in a nonwetted mode.

Figure 15 shows the absorption of pure  $\text{SO}_2$  and of  $\text{SO}_2$  from  $\text{SO}_2$ -air mixtures when the absorbent liquid, water, flows in the shell side. Pure  $\text{SO}_2$  removals as high as 99% were achieved with a liquid flow rate of about  $40 \text{ cm}^3/\text{min}$  in a 30-cm-long, 0.45-cm-diameter module 1 with a gas flow rate of about  $433 \text{ cm}^3/\text{min}$ . It is useful to know in this context that absorption of flue gases can be performed effectively using seawater (Radojevic and Tressider, 1992).

Considerable data were taken also when either pure  $\text{CO}_2$  or  $\text{CO}_2-\text{N}_2$  mixtures flowed on the tube side while water flowed countercurrently in the shell side (Karoor, 1992). The gas absorption rates were always significantly lower in most of the systems studied compared to the tube-side liquid flow cases due most likely to shell-side channeling. A much better gas absorption performance with shell-

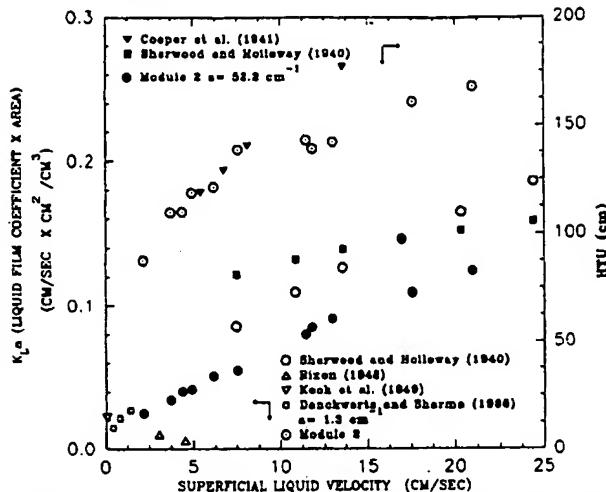


Figure 16. Comparison of  $K_{L\alpha}$  and HTU of a hollow fiber contactor to that of packed towers for  $\text{CO}_2$  absorption in water.

side liquid flow was achieved using a crossflow module (Karoor, 1992). These results are not being reported here.

The pressure drop encountered in the gas-phase flow through the shell side was always very small for pure gas stream used here; mostly it was negligible. In a few cases the  $\Delta P$ s encountered varied between 0.05 and 0.3 psi. When these gases flowed in the tube side, the  $\Delta P$  was usually higher, in a few cases as much as 3–5 psi.

The absorption performance of a hollow fiber module needs to be compared with that of a conventional packed bed absorber. To this end, several experimental observations in a packed bed column for the  $\text{CO}_2-\text{H}_2\text{O}$  system reported in the literature are now compared to those of hollow fiber modules in Figure 16 using the data from the hollow fiber module 2 (dimensions of which are given in Table I). In this plot the equipment-average value of  $K_{L\alpha}$  obtained from the product of experimental  $K_{0lm}$  (eq 1) and  $a$  is plotted as a function of superficial liquid velocity ( $a$  is the interfacial gas-liquid contact area per unit volume). To obtain the superficial liquid velocity in a hollow fiber module in a manner equivalent to that defined for a packed bed, the following definition of superficial velocity is used for a hollow fiber contactor:

$$F_1 = \frac{F_1}{60(\text{empty cross-sectional area})} \quad (30)$$

where  $F_1$  is the liquid flow rate in  $\text{cm}^3/\text{min}$ , empty cross-sectional area is defined as  $\pi r_e^2$ , and  $r_e$  is the shell inner radius. From the figure it is observed that, as the superficial velocity increases, the  $K_{L\alpha}$  also increases. Rixon (1948) studied the absorption of  $\text{CO}_2$  in a commercial size plant where the column inner diameter was 27 in. and the packed height was 18 ft. The packing used was 1 1/2-in.  $\times$  1 1/2-in. stoneware. The  $K_{L\alpha}$  values obtained by them are quite low. Danckwerts and Sharma (1966) studied the absorption of  $\text{CO}_2$  in water in a column of 18-in. diameter and the packing material 1 1/2-in. ceramic raschig rings. The values of  $K_{L\alpha}$  obtained by them were also low when compared to those in the hollow fiber module. In their case, the liquid film coefficients were very high but the area/volume was only 1.3  $\text{cm}^2/\text{cm}^3$ , which resulted in low  $K_{L\alpha}$  values. Sherwood and Holloway (1940) used their  $\text{O}_2$  desorption data and came up with general correlations for calculating  $K_{L\alpha}$  and HTU for packed beds with various packing materials:

$$\frac{K_{L\alpha}}{D_L} = \alpha \left( \frac{L}{\mu_L} \right)^{1-n} \left( \frac{\mu_L}{\rho_L D_L} \right)^{1-s} \quad (31)$$

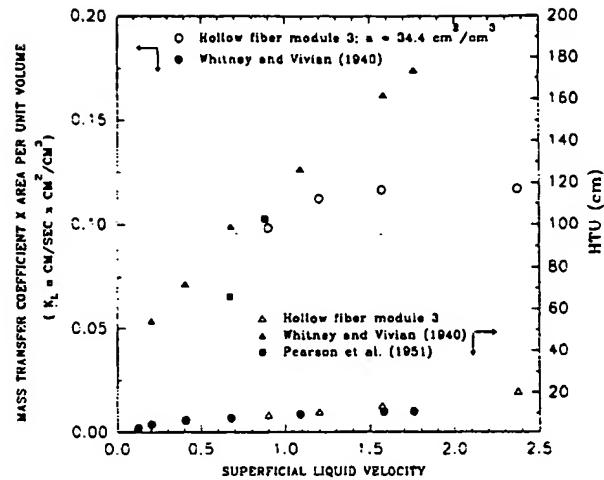


Figure 17. Comparison of  $K_{L\alpha}$  and HTU of a hollow fiber contactor to that of packed towers for absorption of  $\text{SO}_2$  in water.

where  $K_{L\alpha}$  is the mass-transfer coefficient for liquid phase,  $\text{h}^{-1}$ ;  $H_L$  is the height of an individual liquid-film transfer unit, ft;  $L$  is the liquid mass velocity,  $\text{lb}/(\text{ft})^2(\text{h})$ ;  $\mu_L$  is the viscosity of the liquid,  $\text{lb mass}/(\text{ft})(\text{h})$ ;  $\rho_L$  is the density of the liquid,  $\text{lb}/\text{ft}^3$ ;  $D_L$  is the diffusion coefficient for liquid phase,  $\text{ft}^2/\text{h}$  (approximately  $6.6 \times 10^{-5} \text{ ft}^2/\text{h}$  for  $\text{CO}_2$  at 18 °C); and  $\alpha$ ,  $n$ , and  $s$  are constants for various packings used. The above equations provide general estimates of  $K_{L\alpha}$  and HTU values at low liquid flow rates. The above equations can also be used for absorption of  $\text{CO}_2$  as the gas-film resistance is generally neglected. As seen in Figure 16, the  $K_{L\alpha}$  values from such an equation are always smaller than those obtained in the hollow fiber contactor. Also most of the data for packed beds are available for low liquid flow rates; at higher liquid flow rates there may be flooding or loading. Thus, even with a simple absorbent liquid used in these experiments, a remarkably higher amount of removal was achieved with relatively small modules and flow velocities. The potential for increased removal is enormous with the application of reactive absorbents and appropriate operating conditions.

Figure 16 shows also the HTU variation as a function of the superficial liquid velocity. The HTU in the case of a hollow fiber is given as

$$\text{HTU} = \frac{v_{sL}}{K_{L\alpha}} \quad (33)$$

In the figure the data obtained by Cooper et al. (1941) in a 2-in.  $\times$  2 1/16-in. steel raschig ring column show the highest HTU for a given superficial velocity. The HTU given by eq 32 (Sherwood and Holloway, 1940) is also plotted. As seen in the figure, the equation still yields higher HTU when compared to that for a hollow fiber contactor. The difference is especially large at lower superficial velocities.

Figure 17 compares the  $K_{L\alpha}$  values for hollow fiber module 3 to those of packed towers for  $\text{SO}_2$  absorption in water. The  $K_{L\alpha}$  is plotted as a function of the superficial velocity. The  $K_{L\alpha}$  increases as the superficial velocity increases. Whitney and Vivian (1949) studied the absorption of  $\text{SO}_2$  in water in a 8-in.-diameter tower using 1-in. ceramic packing; as observed from the figure the  $K_{L\alpha}$  obtained by them were low. At a superficial velocity of 1.5 cm/s, the  $K_{L\alpha}$  obtained in a hollow fiber module is

around 10 times higher than that obtained in a packed tower. Figure 17 compares also the HTU of hollow fiber module 2 to that of packed towers. The data obtained by Whitney and Vivian (1949) and Pearson et al. (1951) are plotted as a function of the superficial velocity. As observed from the figure, very low HTU values were obtained when a hollow fiber module was used for gas-liquid contacting when compared to the packed towers.

## Conclusions

For CO<sub>2</sub> absorption in water and SO<sub>2</sub> absorption in water in hydrophobic microporous hollow fiber modules, both liquid and gas flow rates were controlled independently, and there were no operational problems in either wetted or nonwetted modes of operation. For CO<sub>2</sub> absorption in water, the wetted mode of operation offers considerably higher resistance to mass transfer when compared to the nonwetted mode of operation. It is a liquid-phase-controlled process; thus filling the membrane pores with water increases the resistance to mass transfer. Increasing the contact area in a given module or the packing fraction for a given fiber size increases the CO<sub>2</sub> removal capacity. The contact area in a hollow fiber device can be increased by increasing the number of fibers without any substantial increase in the physical size of the contacting device. When the absorbent liquid, water, flows in the tube side, simulation results from a model based on first principles are in good agreement with experimental observations. The Leveque solution also agrees well with the experimental results. The experimentally obtained  $K_{L\alpha}$  values for the contactors are considerably larger than those for packed towers. Even though the mass-transfer coefficients obtained for hollow fiber modules in this study are similar to or lower than the values reported for packed beds in the literature, the contact area per unit equipment volume in the case of hollow fiber devices is many times higher than that for packed towers. Therefore, at a superficial liquid velocity of 1.25 cm/s, the  $K_{L\alpha}$  value was 5 times higher than that for a conventional packed tower. As the liquid flow rate increases, the difference in the  $K_{L\alpha}$  values is decreased. At higher flow rates, there are operational problems in packed tower like flooding and loading; there is no such problem in a hollow fiber module as long as the correct phase pressure conditions are maintained. The HTU in a hollow fiber module is several times lower than that in a packed tower operating under similar flow conditions.

For the absorption of SO<sub>2</sub> in pure water in a microporous hollow fiber contactor, the following conclusions are relevant. The  $K_{L\alpha}$  value for a hollow fiber contactor was about 10 times higher than that for conventional packed towers for a superficial liquid velocity of 1.5 cm/s. The HTUs obtained were also about 10 times lower than those obtained for a packed tower for SO<sub>2</sub> absorption in water. Thus efficient contacting is achieved in the case of a hollow fiber when compared to conventional contactors. In the case of SO<sub>2</sub> absorption in water, removal rates as high as 99% were attained in module 1 (a 1/4-in. tube containing 102 X-10 fibers, 30 cm long). Thus, even in a small module high removal rates were attained, proving the efficiency of the contactor. Incorporation of SO<sub>2</sub> dissociation reaction in pure water predicts 10%–15% higher SO<sub>2</sub> concentration in the liquid outlet. The wetted mode of operation with pure water as absorbent offers somewhat higher resistance when compared with the nonwetted mode of operation. Thus it is advantageous to operate the module in the nonwetted mode. Higher contact areas per unit module volume provide higher SO<sub>2</sub> removal efficiencies. The

model simulations are in good agreement with the experimental observations for cases where the absorbent liquid flows in the tube side.

## Acknowledgment

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## Nomenclature

- $a$  = gas-liquid interfacial area/contactor volume (cm<sup>-1</sup>)
- $A_T$  = mass-transfer area based on surface area of gas-liquid contact (cm<sup>2</sup>)
- $C_{ijk}$  = concentration of species  $i$  in the phase  $j$  ( $j = g$ , for gas, and  $j = l$ , for liquid) and in region  $k$  ( $k = m$ , for fiber wall pore,  $k = s$ , for shell side, and  $k = t$ , for fiber bore) (mol/cm<sup>3</sup>)
- $C_{Aw}$  = concentration of species A in liquid at the gas-liquid interface at the membrane wall (mol/cm<sup>3</sup>)
- $C_{Bw}$  = concentration of species B in liquid at the gas-liquid interface at the membrane wall (mol/cm<sup>3</sup>)
- $C_{g\ in}$  = concentration of the absorbed gas species at the gas stream inlet (mol/cm<sup>3</sup>)
- $C_{g\ out}$  = concentration of the absorbed gas species at the gas stream outlet (mol/cm<sup>3</sup>)
- $C_{i\ g}$  = concentration of species  $i$  in the gas phase at the gas-liquid interface (mol/cm<sup>3</sup>)
- $C_{i\ l}$  = concentration of species  $i$  in the liquid phase at the gas-liquid interface (mol/cm<sup>3</sup>)
- $C_{l\ out}$  = concentration of CO<sub>2</sub> or SO<sub>2</sub> in liquid outlet (mol/cm<sup>3</sup>)
- $\Delta C_{lm}$  = logarithmic mean concentration difference of CO<sub>2</sub> or SO<sub>2</sub> (mol/cm<sup>3</sup>)
- $D_{Bl}$  = diffusion coefficient of species B in the liquid phase (cm<sup>2</sup>/s)
- $D_{ik}$  = diffusion coefficient of species  $i$  in region  $k$  ( $k = m$ , for fiber wall pore,  $k = s$ , for shell side, and  $k = t$ , for tube side) (cm<sup>2</sup>/sec)
- $F_{g\ in}$  = gas flow rate at module inlet (cm<sup>3</sup>/min)
- $F_{g\ out}$  = gas flow rate at module outlet (cm<sup>3</sup>/min)
- $H$  = dimensionless Henry's constant as defined by eq 3
- $k_1$  = forward rate constant for SO<sub>2</sub> ionization (s<sup>-1</sup>)
- $K_E$  = equilibrium constant of a chemical reaction
- $K_{Olm}$ ,  $K_L$  = overall mass-transfer coefficient (cm/s)
- $L$  = effective length of the module, cm
- $N_f$  = total number of fibers in the hollow fiber module
- $N_G$  = Graetz number,  $\pi r_t N_R N_S / 2L$
- $N_R$  = Reynolds number,  $2r_t v_m \rho / \mu$
- $N_S$  = Schmidt number,  $\mu / \rho D_A$
- $N_{Sh}$  = Sherwood number,  $K_{Olm} (2r_t) / D_A$
- $P_g$ ,  $P_{eq}$  = gas-phase pressure (psig)
- $P_{eq}$  = liquid-phase pressure (psig)
- $r$  = radial coordinate position (cm)
- $r_o$  = free surface radius (see eq 6) (cm)
- $r_i$  = molar production rate of species  $i$  by reaction per unit volume (mol/(cm<sup>3</sup>·s))
- $r_o$ ,  $r_t$  = outer radius and inner radius of the fiber (cm)
- $r_p$ ,  $r_s$  = pore diameter, shell radius (cm)
- $R_{gas}$  = gas constant (atm·cm<sup>3</sup>/(mol K))
- $T_R$  = room temperature (°C)
- $v_s$ ,  $v_t$  = shell side and tube side velocity (cm/s)
- $V_l$  = liquid flow rate (cm<sup>3</sup>/min)

## Greek Letters

- $\epsilon$  = packing fraction
- $\mu_j$  = viscosity of the fluid in the shell side (gm/(cm·s))
- $\lambda$  = mean free path (cm)
- $\rho$  = density (g/cm<sup>3</sup>)

$\tau$  = tortuosity of the membrane pore

### Subscripts

A, B = species CO<sub>2</sub> and HSO<sub>3</sub><sup>-</sup>, respectively

i = component involved

j = phase involved (j = g, gas, or j = l, liquid)

k = location of the species (k = m, for fiber wall pore, k = s, for shell side, and k = t for tube side)

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